## SEEP2022 Program Summer School Helium cryogenics for superconducting cooling

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#### **Cryoboat Group**

Institute of Refrigeration and Cryogenics College of Energy Engineering, Zhejiang University

#### **Research topics:**

Liquid hydrogen Air separation Cryocooler Helium cryogenics



attracteo **Group Leader:** Prof. Limin Qiu group **Researchers:** Prof. Kai Wang Prof. Xiaoqin Zhi Prof. Shiran Bao Dr. Song Fang 4 **Postdoc:** Dr. Shaolong Zhu Dr. Hanying Jiang

Cryoboat

**Over 20 graduate students** 



He II and Quantum hydrodynamics

Flow visualization studies in He II

Defects location on SRF cavities by MTV

Vacuum loss in LHe cooled tube

**Recent and Future development** 



## He II and Quantum hydrodynamics



#### Specificity of helium and superfluid helium (He II)

- Lowest melting and boiling points of all the elements.
- Helium does not solidify at atmospheric pressure.



#### Peter Kapitza, Nobel Prize in Physics 1978







#### Superfluid helium (He II) as a coolant



Lev Landau, Nobel Prize in Physics 1962

There exist two components in He II:

- Superfluid component (condensate)
- Normal-fluid component (excitations)

	Density	Velocity	Viscosity	Entropy
Superfluid	$\rho_s(T)$	$\mathbf{v}_{s}(\mathbf{r})$	0	0
Normal fluid	$\rho_n(T)$	$\mathbf{v}_n(\mathbf{r})$	$\eta_n(T)$	$s_n(T)$





Two sound modes:

- 1<sup>st</sup> sound: two fluids move in phase
- 2<sup>nd</sup> sound: two fluids move out of phase



Superfluid turbulence is a random tangle of quantized vortices.



#### **Quantum turbulence**

Each vortex has the same conserved circulation. QT is thus simpler to model than classical turbulence.

 Turbulence in normal fluid can be affected by the quantized vortices.



#### **Classical turbulence**

Quasi-particles can interact with vortices  $\rightarrow$  origin of many peculiar hydrodynamics in the two-fluid system.



Heat transfer in He-II is by counterflow :

the superfluid moves towards the source of heat;

the normal fluid flows in the opposite direction, carrying thermal energy.



Turbulence in counterflow can affect the heat transfer !





At low T when there is no normal fluid, the decay of QT energy at scales smaller than vortex spacing is believed to be via Kelvin wave cascade.

Energy RICHARDSON CASCADE 0000000000000 000000000000000000 RECONNECTIONS KELVIN-WAVE CASCADE Thermal excitations

Figure 4. How quantum turbulence might evolve in a superfluid at very low temperature.15 (a) First, turbulent energy mechanically injected at large length scale flows to ever smaller length scales in a quasiclassical Richardson cascade, down to length scale l, the typical distance between quantized vortex lines. (b) Then vortex-line reconnection comes into play. The reconnection is modestly dissipative, producing some phonon emission. But more of the energy goes into generating Kelvin waves of vortices (c). Strong nonlinear interactions make the Kelvin waves cascade down to wavelengths short enough so that their energy can be dissipated as phonons or other thermal excitations. Aspects of this picture are still controversial.

(1) Classical turbulent energy spectrum is indicated but no direct determination.

$$dE/dt = -\nu\omega^2 = -\nu_{eff}(\kappa L)^2$$

(2) What is the spectrum in the Kelvin wave cascade regime ?

(3) Is there any bottleneck at the transition from classical cascade to Kelvin wave cascade ?

Direct vortex-line visualization at low T is urgently needed!



#### Do we really have a complete understanding of counterflow?



Challenge: many tools developed for classical fluids are not applicable We need tools that allow independent quantitative flow measurements in the two-fluid system!



## Flow visualization studies in He II

Particle imaging velocimetry (PIV) with polymer microspheres, solidified hydrogen ice particles



Particle tracking velocimetry (PTV, with hydrogen isotopes ice particles)



Direct visualization of quantized vortices: Particles are observed to bind on vortex lines due to Bernoulli's effect



Bewley, Lathrop, and Sreenivasan Nature 441, 588 (2006)

• Our recent progress with PTV in He II:

#### 1) Thermal counterflow:



(a)

3

0

-2

-3

3

-2 -3

> -7 -6

(c)

-7 -6

(b)

-5

-5

-5

-4 -3

-4

-3 -2

-3 -2 -1 0

-2

-1 0

0

Imaging Plane Width (mm)

maging Plane Height (mm)

maging Plane Height (mm)

maging Plane Height (mm)

-2 -3

> -7 -6

- Particle tracks and V<sub>v</sub> pdfs show two distinct regimes
- Two peaks at low heat fluxes: G1 (lower velocity peak) and G2 (high velocity peak)
- One peak at high heat fluxes: G3



T= 1.85 K

G1 Fit

G2 Fit

10

V<sub>p,m</sub> G1 Fit G2 Fit 15

0 V<sub>p,m</sub>

(d)

0

5

v<sub>p.m</sub> (mm/s)

0

600

400

200

600

400

-5

(e)

Pr(v<sub>p,m</sub>)/∆v (s/mm)

6

5

4

7

3 4 5

2 3

2 3 4 5

1

Imaging Plane Width (mm)

1

Imaging Plane Width (mm)

#### **PIV/PTV techniques using micron-sized tracers**

Track length of the G2 particles  $\rightarrow$  vortex spacing

Track length of G2 can be used to determine the mean free path, s, of the particles through the vortex tangle:



0

-1

Result agrees very well with 2<sup>nd</sup> sound measurement



B. Mastracci and W. Guo, submitting to Phys. Rev. (2018)



#### 2) PTV in grid turbulence:



B. Mastracci and W. Guo, Rev. Sci. Instrum., 89, 015107 (2018)

Particles are expected to trace the coupled flow in grid turbulence.



#### **PIV/PTV techniques using micron-sized tracers**









B. Mastracci and W. Guo, to be submitted to Phys. Rev. (2018)

Some issues of PIV/PTV method:

1. Particles have a wide size distribution and non-spherical shapes; they are not neutrally buoyant; they can form large clusters.

2. Particles interact with both the normal fluid and the vortices. Tracer behavior is hard to interpret.



(D. Zmeev, et al, Phys. Rev. Lett., 110, 175303 (2013))

Metastable He<sub>2</sub><sup>\*</sup> molecules can be easily produced as a result of ionization or excitation in LHe4:

```
e^- + He<sup>+</sup> + He \rightarrow He<sup>*</sup> + He \rightarrow He<sup>*</sup><sub>2</sub>
```

singlet state  $a^{1}\Sigma_{u}^{+}$  lifetime: ~1ns

triplet state  $a^{3}\Sigma_{u}^{+}$  lifetime: ~13s

- Above 1K : molecules trace the normal-fluid component only.
- Below 0.5 K : molecules can be trapped on vortex lines



#### **Imaging He<sub>2</sub>**<sup>\*</sup> **molecules: Laser-induced fluorescence**



W.G. Rellergert et al., Phys. Rev. Lett, 100 (2008).

For molecules in the triplet ground state a(0):

- A 905 nm pulsed laser is used to drive a cycling transition.
- Fluorescent light emitted at 640 nm.





Guo, et al., Phys. Rev. Lett., 102, 235301 (2009)



Guo, et al., Phys. Rev. Lett. 105, 045301 (2010).

#### **Imaging He<sub>2</sub>**<sup>\*</sup> **molecules: Laser-induced fluorescence**

Femtosecond laser field ionization in helium:



 $I \ge 10^{13} \, {\rm W/cm^2}$ 



Pulse length: 35 fs

Pulse energy: up to 4 mJ

Rep rate: up to 5 kHz



J. Gao, et al., Rev. Sci. Instrum. 86, 093904 (2015)



W. Guo, et al., PNAS, 111, 4653 (2014)

- Thin tracer lines can be produced and tracked, allowing high precision flow field measurement.
- This technique is applicable to He II, He-I, and gaseous helium.

Three distinct velocity profiles of the normal fluid were observed.



A. Marakov, G. Jian, et al., Phys. Rev. B 91, 094503 (2015).

 Recent simulations suggest that the mutual friction near the wall may flatten the profile of the normal fluid.

A. W. Baggaley and S. Laizet, Phys. Fluids 25, 115101 (2013)



The velocity PDF in turbulent normal fluid is found to be a Gaussian. The turbulence intensity is higher than classical channel flow:



The 2nd order transverse structure function revealed novel form of turbulence



Shiran BAO – Helium cryogenics for superconducting cooling



## **Defects location on SRF cavities by MTV**

#### Superconducting radio-frequency (SRF) cavities



SRF Cavities Applications:

High Energy Physics: LHC, CEBAF; Radiation Sources: XFEL, ERL; Nuclear Physics: ATLAS, TRIUMF, SNS, EURSOL; Upcoming: ESS, FRIB, PIP-II ...

- SRF cavities are key components in many modern particle accelerators due to their high Q factors.
- The maximum acceleration gradient is limited by the 'quench' of SRF cavities.

T. Khabiboulline, Engineering for Particle Accelerators, (2017).



3

#### Superconducting RF cavities and surface quench spots



- The quenching spot causes transient heat transfer into He II (≈1 ms, 1-10 J).
- The maximum achievable field can be improved by removing the surface defects by mechanical grinding, tumbling the cavity, and electron or laser re-melting.
- Precise location of the hotspots is prerequisite for a successful repair.



Typical defects images captured by Cornell University and DESY



#### Surface T-mapping



 Installation and maintenance required for over 1000 temperature sensors and wires.

#### **Rotational T-mapping**



 Conducted just below the quench threshold. At least 2 cool downs and reconfiguration of the sensor are normally needed for good accuracy



#### **Existing methods: (2). 2<sup>nd</sup> sound trilateration**





Conway, et al., Supercond. Sci. Technol. 30, 034002 (2017).



- Hotspot leads to emission of 2nd sound waves in He II.
- Converged location can have an uncertainty of 5-10 mm.
- Convergence of the signals requires a 2nd sound speed different than known value.



#### Possible non-contacting detection via flow visualization



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#### **Proof-of-concept experiment using SMD heaters**





- An array of thin film resistors (50 Ohm, 0.8x0.8 mm<sup>2</sup>).
- The PCB is in He II at controlled temperatures (1.85 K).
- A voltage pulse (1-8 ms) is applied to a chosen heater (78-287 W/cm<sup>2</sup>).

S. Bao and W. Guo, arXiv:1812.07080 (2018) (submitting to Phys. Rev. Applied)





#### Typical baseline and deformed tracer-line images



$$dr = v_n \cdot dt = \frac{q(r)}{\rho sT} \cdot dt \qquad r_f^3 = r_0^3 + \frac{3}{2\pi\rho sT} \int_{t_0}^{t_0 + \Delta t} \dot{Q}_s dt = r_0^3 + \frac{3Q_s}{2\pi\rho sT}$$

#### Two fitting parameters:

1) The location of the heater:  $x_0$  2) The total energy carried by second-sound:  $Q_s$ 



 The obtained x<sub>0</sub> is always within a few hundred microns from the actual heater location, clearly proves the feasibility of this novel technology.







#### Possible origin of the "fast" second sound



 A new explanation for the decadeslong puzzle observed in previous second-sound triangulation experiments is proposed.  The size of this cavitation zone is estimated based on the knowledge obtained about the transported heat. (15 W/cm<sup>2</sup> threshold is selected)



#### Our long-term goal:

#### **Rotational quick scan of real cavities**



**3D imaging of 2 tracer-lines** 



#### 3D imaging: Stereoscopic MTV experimental setup

 We built up a stereoscopic MTV system to capture the 3D profile of the tracer line for detecting a hot spot on a 2D surface.



S.R. Bao and W. Guo, Int. J. Heat Mass Transf., 161, 120259 (2020).





#### **3D imaging: Reconstruction of the tracer-line profile**





- Tracer line is created at about 4-5 mm away from the heater center.
- Reconstruction is performed based on the marker positions.



In the current experiment the heater size is increased by nearly nine times in order to better represent a quench spot, which makes the point heatsource model inapplicable.



We divide the assumed heater into n×m small elements and treat each element as a standalone point heat source. The velocity of a tracer-line segment located at r:

$$\boldsymbol{v}(\boldsymbol{r},t) = \sum_{i=1}^{n} \sum_{j=1}^{m} \boldsymbol{v}_{i,j}(\boldsymbol{r},t) = \sum_{i=1}^{n} \sum_{j=1}^{m} \frac{\boldsymbol{q}_{i,j}(t)}{\rho sT}$$

$$q_{i,j}(t) = \dot{Q}_{i,j}(t)/2\pi R_{i,j}^2$$
  $R_{i,j} = |\mathbf{r}_{i,j} - \mathbf{r}|$ 

• The heat transfer rate is given by:

$$\dot{Q}_{i,j}(t) = \begin{cases} 0 & c_2 t < R_{i,j} \\ \frac{Q_s / \Delta t}{n \times m} & R_{i,j} \le c_2 t \le R_{i,j} + c_2 \Delta t \\ 0 & c_2 t > R_{i,j} + c_2 \Delta t \end{cases}$$





 Our analysis nicely reproduces the heater center with an uncertainty of a few hundred microns (a few percent of the heater size), regardless of the applied heat fluxes.



## Vacuum loss in LHe cooled tube

#### Particle accelerators composed of LHe cooled cryomodules

![](_page_39_Picture_1.jpeg)

![](_page_40_Picture_0.jpeg)

A sudden catastrophic loss of vacuum is one of the most serious failures in LHe cooled particle-accelerator systems

Insulation vacuum space

(usually isolated per cryomodule)

#### Vacuum beamline

(interconnected to create a beamline up to kilometers)

- Loss of vacuum insulation leads to explosive boiloff
  - Equipment damage
  - Personnel injury
  - Oxygen deficiency
  - Gas propagating in the beam tube could affect the entire system
  - Introduce contamination to the inner surface of cavities

![](_page_40_Picture_12.jpeg)

- Example: LHC, CERN 2008
- Beam tube vacuum failure
- 53 cryomodules contaminated/damaged
- 6 tons helium lost
- 6 months to repair

Developing a clear understanding of the complex dynamical heat and mass transfer processes involved following a sudden vacuum break is of great importance for the safe operation of accelerators.

![](_page_41_Picture_0.jpeg)

• Early experiments conducted in straight tube immersed in He I by Dhuley and Van Sciver in our lab confirmed the slowing down propagation

![](_page_41_Figure_3.jpeg)

![](_page_41_Picture_4.jpeg)

![](_page_41_Figure_5.jpeg)

![](_page_41_Figure_6.jpeg)

![](_page_42_Picture_0.jpeg)

Dhuley and Van Sciver attributed the exponential slowing of the gas front propagation to gas condensing and freezing to the walls. They proposed a simple model based on conservation of mass analysis in one dimensional tube flow.

![](_page_42_Figure_3.jpeg)

*Int J Heat Mass Transf* 2016;96:573–81; *Int. J. Heat Mass Transf* 2016;98:728–37

![](_page_43_Picture_0.jpeg)

#### Modified experimental setup using a helical tube system

![](_page_43_Figure_2.jpeg)

eleleleletstatetelele

![](_page_44_Picture_0.jpeg)

#### Typical result for gas propagation in He I and He II cooled tube

![](_page_44_Figure_2.jpeg)

![](_page_45_Picture_0.jpeg)

A 1-D finite deference model that systematically describes the fluid flow, heat transfer and mass deposition of a propagating and condensing gas inside a liquid helium cooled tube has been established.

![](_page_45_Figure_3.jpeg)

- 1. GN2 is treated as an ideal gas:  $PM_{\rm g} = \rho_{\rm g} RT_{\rm g}$  viscosity and gravity are neglected ;
- 2. An experimentally measured initial temperature profile is used for the insulated inlet section;
- 3. Steady-state heat transfer in LHe as the onset time of film boiling is very short;
- 4. The inlet mass flow rate is determined based on the measured pressure change in the tank;
- 5. The GN2 flow velocity assumed to be at the local speed of sound at the inlet.

• The equations that describe the gas propagation in a long pipe cooled by helium:

Mass conservation:

n: 
$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) = -\frac{4}{D_1}\dot{m}_c$$

Momentum conservation:

$$\frac{\partial}{\partial t}(\rho u) + \frac{\partial}{\partial x}(\rho u^2) = -\frac{\partial P}{\partial x} - \frac{4}{D_1}\dot{m}_c u$$

Energy conservation:

$$\frac{\partial}{\partial t} \left[ \rho \left( \varepsilon + \frac{1}{2} u^2 \right) \right] + \frac{\partial}{\partial x} \left[ \rho u \left( \varepsilon + \frac{1}{2} u^2 + \frac{P}{\rho} \right) \right] = -\frac{4}{D_1} \dot{m}_c \left( \varepsilon + \frac{1}{2} u^2 + \frac{P}{\rho} \right) - \frac{4}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu} \cdot k \left( T_g - T_s \right) + \frac{1}{D_1^2} \operatorname{Nu$$

Mass deposition rate: 
$$\dot{m}_c = \sqrt{\frac{M_g}{2\pi R}} \left( \Gamma \sigma_c \frac{P}{\sqrt{T_g}} - \sigma_e \frac{P_w}{\sqrt{T_w}} \right)$$

Radial heat transfer:

$$\rho_w C_w \frac{D_2^2 - D_1^2}{4D_1} \frac{\partial T_s}{\partial t} = q - q_{He} \frac{D_2}{D_1}$$

$$q = \dot{m}_c \left[ \frac{1}{2} u^2 + \hat{h} (T_g, P) - \hat{h} (T_s, P) \right] - \frac{\mathrm{Nu} \cdot k}{D_1} (T_g - T_s)$$

Simplified radial heat transfer model:

![](_page_47_Figure_2.jpeg)

Detailed information about the numerical model can be found in: *Int. Heat Mass Transf.* 129 (2019) 1144–1150; *Int. Heat Mass Transf.* 146 (2020), 118883; *Int. Heat Mass Transf.* 181 (2021) 121885.

![](_page_48_Figure_1.jpeg)

![](_page_49_Picture_0.jpeg)

We see a good agreement for the results with different mass flow rate with a single value of  $B_w$  parameter.

![](_page_49_Figure_2.jpeg)

Threshold Level 4.7 K

![](_page_50_Picture_0.jpeg)

 The heat transfer condition varies greatly with different mass flow rate. We introduced an other tuning parameter ψ to the peak heat flux model which represents the derivation of the situation from an ideal cylindrical geometry.

Tank Pressure (kPa)	Optimal $\psi$	
50	0.43	
100	0.59	
150	1.97	
200	1.96	

![](_page_50_Figure_4.jpeg)

![](_page_51_Picture_0.jpeg)

- The freeze-out length is the upper limit of the propagation of possible contaminant, which can be evaluated with our 1D numerical model
- By analyzing the arrival time of gas front at each position, the freeze-out point can be defined by a linear fitting

![](_page_51_Figure_4.jpeg)

![](_page_52_Picture_0.jpeg)

#### **Empirical formula for the freeze-out length**

For He I cooled tube

![](_page_52_Figure_3.jpeg)

![](_page_52_Figure_4.jpeg)

For He II cooled tube (*T*=1.9 K, *h*=50 cm,  $\psi$ =1)

Empirical	Parameter	Optimal value for He I	Optimal value for He II	Unit
Formula:	а	0.074077	0.01786765	m <sup>2c-b+1</sup> kg <sup>-c</sup> s <sup>c</sup>
$x_{E} = a D_{i}^{b} \dot{m}^{c}$	b	0.91460513	1.02281107	1
$n_F = \alpha D_1 m$	С	1.08364073	1.39491514	1

![](_page_53_Picture_0.jpeg)

## **Recent and Future development**

![](_page_54_Picture_0.jpeg)

Liquid helium-4 has extremely small kinematic viscosity:

![](_page_54_Picture_3.jpeg)

![](_page_54_Picture_4.jpeg)

![](_page_54_Figure_5.jpeg)

![](_page_54_Picture_6.jpeg)

 Channel flows with Re~10<sup>7</sup>
has already been achieved in our cryogenics lab in He-II.

![](_page_55_Picture_0.jpeg)

#### **1. High turbulence test facility in NHMFL**

 Additional fs-laser beam path has been developed to created 2 tracer lines with adjustable distance, this upgrade will enable the 2D full field flow visualization.

![](_page_55_Picture_3.jpeg)

![](_page_55_Picture_4.jpeg)

 Thermal neutron absorption on He3 atoms leads to the production of small clusters of helium molecules.

![](_page_56_Figure_2.jpeg)

#### Cloud density can be tuned by changing He3 concentration or neutron flux.

![](_page_56_Figure_4.jpeg)

![](_page_57_Picture_0.jpeg)

#### 2. PIV/PTV with He<sub>2</sub><sup>\*</sup> clouds created by He3-neutron absorption

 Collaboration with researchers at Oak Ridge and Univ. Tennessee:

Prof. M. R. Fitzsimmons, Prof. Xin Tong, et al

![](_page_57_Picture_4.jpeg)

![](_page_57_Picture_5.jpeg)

![](_page_57_Picture_6.jpeg)

 He<sub>2</sub>\* clouds created by He3neutron absorption has been successfully visualized, 1st report has been published on PRL.

#### 2. PIV/PTV with He<sub>2</sub><sup>\*</sup> clouds created by He3-neutron absorption

10000

![](_page_58_Figure_1.jpeg)

10<sup>4</sup>

105

x-direction z-direction

28

30

#### Reference

- [1] N. Garceau, S.R. Bao, W. Guo, Heat and mass transfer during a sudden loss of vacuum in a liquid helium cooled tube Part III: Heat deposition in He II, International Journal of Heat and Mass Transfer. 181 (2021) 121885. https://doi.org/10.1016/j.ijheatmasstransfer.2021.121885.
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# Thank you !

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